Cyclic Central Configurations in the Four-Body Problem

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A *central configuration* (c.c.) is a configuration of bodies $(\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_n), \mathbf{x}_i \in \mathbb{R}^d$ such that the acceleration vector for each body is a common scalar multiple of its position vector. Specifically, in the Newtonian n-body problem with the center of mass at the origin, for each index i,

$$\sum_{j\neq i}^{n} \frac{m_i m_j(\mathbf{x}_j - \mathbf{x}_i)}{||\mathbf{x}_j - \mathbf{x}_i||^3} + \lambda m_i \mathbf{x}_i = 0$$

for some scalar λ .

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- The collinear c.c.'s correspond to d = 1, planar c.c.'s to d = 2, spatial c.c.'s to d = 3. One can also study theoretically the case d > 3.
- Summing together the n equations above quickly yields $\sum_{i} m_{i} \mathbf{x}_{i} = 0$.



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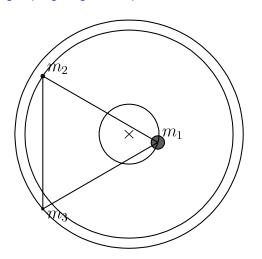
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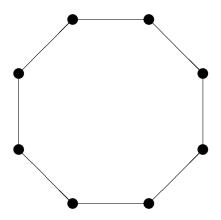
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- 209 articles found on MathSciNet using a general search for "central configurations"

Equilateral Triangle (Lagrange 1772)



Regular *n*-gon (equal mass required for $n \ge 4$)



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- If the center of mass coincides with the center of the circle, the only possibility is the square with equal masses (Hampton, 2003). If such a co-circular c.c. existed for more bodies, other than the regular *n*-gon, this would be a choreography (all bodies tracing out the same curve) without equal time spacing between bodies.

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- Spatial 5-body pyramidal c.c.'s exist where four bodies lie on a sphere with the fifth body at the center of the sphere. The four bodies forming the base of the configuration are co-circular and form a c.c.c. in the 4-body problem (Fayçal, 1996; Albouy, 2003)

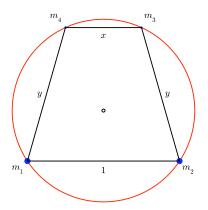


Figure: An example of a cyclic central configuration. The center of the circumscribing circle is marked with an O.

An Alternate Characterization of CC's

Let $r_{ij} = ||\mathbf{q}_i - \mathbf{q}_j||$ where \mathbf{q}_i denotes the position of the *i*-th body. The Newtonian potential function is

$$U(\mathbf{q}) = \sum_{i < j}^{n} \frac{m_i m_j}{r_{ij}}$$

The equations of motion for the *n*-body problem are then given by

$$m_i\ddot{\mathbf{q}}_i = \frac{\partial U}{\partial \mathbf{q}_i}, \quad i \in \{1, 2, \dots n\}$$

$$= \sum_{i \neq i}^n \frac{m_i m_j (\mathbf{q}_j - \mathbf{q}_i)}{r_{ij}^3}$$

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Consequently, the *i*-th equation for a c.c. can be written as

$$\frac{\partial U}{\partial \mathbf{q}_i}(\mathbf{x}) + \lambda m_i \mathbf{x}_i = 0.$$



CC's as critical points of U

The moment of inertia $I(\mathbf{q})$ (w.r.t. the center of mass) is defined as

$$I(\mathbf{q}) = \frac{1}{2} \sum_{i=1}^{n} m_i ||\mathbf{q}_i||^2.$$

Thus, the equations for a c.c. can be viewed as a Lagrange multiplier problem (set $I(\mathbf{q}) = k$):

$$\nabla U(\mathbf{x}) + \lambda \, \nabla I(\mathbf{x}) = 0$$

where $\mathbf{x} = (\mathbf{x}_1, \dots \mathbf{x}_n)$.

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In other words, a c.c. is a critical point of U subject to the constraint I=k (the mass ellipsoid). This gives a useful topological approach to studying central configurations (Smale, Conley, Meyer, McCord, etc.)

Mutual Distances Make Great Coordinates

Recall:

$$U(\mathbf{q}) = \sum_{i < j}^{4} \frac{m_i m_j}{r_{ij}}$$

Alternative formula for *I* in terms of mutual distances: (center of mass at origin)

$$I(\mathbf{q}) = \frac{1}{2M} \sum_{i < j}^4 m_i m_j r_{ij}^2$$

where $M = m_1 + \cdots + m_4$ is the total mass.

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Problem: The six variables r_{12} , r_{13} , r_{14} , r_{23} , r_{24} and r_{34} are not independent in the planar problem. Generically, they describe a tetrahedron, not a planar configuration. It is easy to see that the regular tetrahedron is the only non-planar c.c. in the four-body problem.

4-body Planar CC's

To use the six mutual distances r_{12} , r_{13} , r_{14} , r_{23} , r_{24} , r_{34} as variables, we need an additional constraint that ensures the configuration is planar. We require that the volume of the tetrahedron be zero (Cayley-Menger determinant).

$$F = \begin{vmatrix} 0 & 1 & 1 & 1 & 1 \\ 1 & 0 & r_{12}^2 & r_{13}^2 & r_{14}^2 \\ 1 & r_{12}^2 & 0 & r_{23}^2 & r_{24}^2 \\ 1 & r_{13}^2 & r_{23}^2 & 0 & r_{34}^2 \\ 1 & r_{14}^2 & r_{24}^2 & r_{34}^2 & 0 \end{vmatrix} = 0$$

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Cumbersome! Is there a simpler constraint that ensures both planarity and that the bodies lie on a circle?

Ptolemy's Theorem

If four bodies lie on a common circle and are numbered sequentially (ie. the diagonals have lengths r_{13} and r_{24}), then P=0, where

$$P = r_{12}r_{34} + r_{14}r_{23} - r_{13}r_{24}.$$

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Theorem (Apostol, 1967): For any convex quadrilateral numbered sequentially or for any tetrahedron, $P \ge 0$ with equality iff the four bodies lie on a circle.

If we restrict to the space of geometrically realizable configurations, $r_{ij} > 0$ and $r_{ij} + r_{jk} > r_{ik}$ for all possible triples of indices (i, j, k), then P = 0 iff the configuration is cyclic (with a sequential ordering.)

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Goal: Find critical points of $U + \lambda M(I - I_0) + \sigma P$ satisfying $I = I_0$ and P = 0.

Using the six mutual distances as variables, we find

$$m_{1}m_{2}(r_{12}^{-3}-\lambda)=\sigma\frac{r_{34}}{r_{12}}, \qquad m_{3}m_{4}(r_{34}^{-3}-\lambda)=\sigma\frac{r_{12}}{r_{34}}$$

$$m_{1}m_{3}(r_{13}^{-3}-\lambda)=-\sigma\frac{r_{24}}{r_{13}}, \qquad m_{2}m_{4}(r_{24}^{-3}-\lambda)=-\sigma\frac{r_{13}}{r_{24}}$$

$$m_{1}m_{4}(r_{14}^{-3}-\lambda)=\sigma\frac{r_{23}}{r_{14}}, \qquad m_{2}m_{3}(r_{23}^{-3}-\lambda)=\sigma\frac{r_{14}}{r_{23}}.$$

Using the six mutual distances as variables, we find

$$m_1 m_2 (r_{12}^{-3} - \lambda) = \sigma \frac{r_{34}}{r_{12}}, \qquad m_3 m_4 (r_{34}^{-3} - \lambda) = \sigma \frac{r_{12}}{r_{34}}$$

$$m_1 m_3 (r_{13}^{-3} - \lambda) = -\sigma \frac{r_{24}}{r_{13}}, \qquad m_2 m_4 (r_{24}^{-3} - \lambda) = -\sigma \frac{r_{13}}{r_{24}}$$

$$m_1 m_4 (r_{14}^{-3} - \lambda) = \sigma \frac{r_{23}}{r_{14}}, \qquad m_2 m_3 (r_{23}^{-3} - \lambda) = \sigma \frac{r_{14}}{r_{23}}.$$

This yields a well-known relation of Dziobek (1900)

$$(r_{12}^{-3} - \lambda)(r_{34}^{-3} - \lambda) = (r_{13}^{-3} - \lambda)(r_{24}^{-3} - \lambda) = (r_{14}^{-3} - \lambda)(r_{23}^{-3} - \lambda)$$
 (1)

which must be true for any planar 4-body c.c., not just cyclic c.c.'s.

Eliminating λ from equation (1) in a clever way yields

$$(r_{13}^3 - r_{12}^3)(r_{23}^3 - r_{34}^3)(r_{24}^3 - r_{14}^3) = (r_{12}^3 - r_{14}^3)(r_{24}^3 - r_{34}^3)(r_{13}^3 - r_{23}^3).$$
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Equation (2) is necessary and sufficient for a 4-body planar c.c. given that the six mutual distances determine a geometrically realizable planar configuration. However, it does *not* ensure positivity of the masses.

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Generically, the space of 4-body planar c.c.'s is three dimensional since it is described by the equations F = 0, $I = I_0$ and equation (2). Restricting to a circle yields a two dimensional space.

Mass Ratios

$$\frac{m_2}{m_1} = \frac{(\lambda - r_{13}^{-3}) r_{13} r_{14}}{(r_{23}^{-3} - \lambda) r_{23} r_{24}} = \frac{(r_{14}^{-3} - \lambda) r_{13} r_{14}}{(\lambda - r_{24}^{-3}) r_{23} r_{24}}$$

$$\frac{m_3}{m_1} = \frac{(r_{12}^{-3} - \lambda) r_{12} r_{14}}{(r_{23}^{-3} - \lambda) r_{23} r_{34}} = \frac{(r_{14}^{-3} - \lambda) r_{12} r_{14}}{(r_{34}^{-3} - \lambda) r_{23} r_{34}}$$

$$\frac{m_4}{m_1} = \frac{(r_{12}^{-3} - \lambda) r_{12} r_{13}}{(\lambda - r_{24}^{-3}) r_{24} r_{34}} = \frac{(\lambda - r_{13}^{-3}) r_{12} r_{13}}{(r_{34}^{-3} - \lambda) r_{24} r_{34}}.$$

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Without loss of generality, let $r_{12} = 1$ be the longest exterior side and let $r_{14} \ge r_{23}$. Requiring positive masses gives

$$r_{13} \geq r_{24} > r_{12} = 1 \geq r_{14} \geq r_{23} \geq r_{34}.$$
 (3)

The diagonals are longer than any of the exterior sides. The longest exterior side is opposite the smallest.

Symmetric Example: Kite Configurations

A quadrilateral is a kite if two opposite bodies lie on an axis of symmetry of the configuration. Equation (2) is immediately satisfied when $r_{12} = r_{14} = 1$ and $r_{23} = r_{34}$, yielding a kite configuration with an axis of symmetry between bodies 1 and 3.

Set $r_{23} = r_{34} = x$ and the diagonals $r_{13} = c$ and $r_{24} = 2x/c$ where $x \in (0, 1]$ is a parameter and $c = \sqrt{1 + x^2}$. These distances satisfy Ptolemy's relation P = 0.

We find $m_2=m_4$, as expected from symmetry. Setting $m_1=1$, we have $m_2=m_4=m$ and $m_3=\alpha m$ where

$$m = \frac{4x(c^3 - 1)}{c(8 - c^3)}$$
 and $\alpha = \frac{c(8x^3 - c^3)}{4(c^3 - x^3)}$.

To ensure positive masses, we must have $1/\sqrt{3} < x \le 1$.

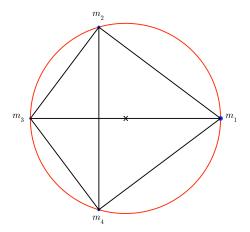


Figure: An example of a cyclic kite central configuration with $m_2 = m_4$.

Theorem

There exists a one-parameter family of cyclic kite central configurations with bodies one and three lying on the diameter of the circumscribing circle. The masses are $m_1 = 1$, $m_2 = m_4 = m$ and $m_3 = \alpha$ m and are ordered $m_1 \geq m_2 = m_4 \geq m_3$ with equality iff the configuration is a square. At one end of the family $(x = 1/\sqrt{3})$ is a c.c. of the planar, restricted 4-body problem, with bodies 1, 2 and 4 forming an equilateral triangle and $m_3 = 0$. At the other end (x = 1) is the square with equal masses.

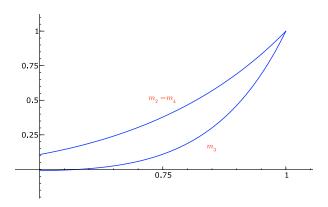


Figure: The values of the masses for the cyclic kite c.c.'s

The longest arc along the cirmcumscribing circle, denoted θ_{12} , occurs between bodies 1 and 2. We have that

$$\theta_{12} = \pi - 2 \arctan(x)$$

and the maximum arc is 120° ($m_3 = 0$) while the minimum arc is 90° (square).

Theorem

The supremum of the largest interior angle of the kite ccc is 120° while the infimum of the smallest interior angle is 60°. Moreover, the arc length along the circumscribing circle between bodies 1 and 2 is a decreasing function of the parameter x with a supremum of $2\pi/(3\sqrt{3})\approx 1.2092$ and a minimum attained at the square configuration of $2^{-3/2}\pi\approx 1.1107$.

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A cyclic central configuration is an isosceles trapezoid if and only if $m_1 = m_2$ and $m_3 = m_4$.

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Key calculation in proof: If $m_1 = m_2$, then

$$r_{23}^2 r_{24}^2 (r_{13}^3 - r_{14}^3) - r_{13}^2 r_{14}^2 (r_{24}^3 - r_{23}^3) = 0.$$
(4)

Likewise, if $m_3 = m_4$, then

$$r_{13}^2 r_{23}^2 (r_{24}^3 - r_{14}^3) - r_{14}^2 r_{24}^2 (r_{13}^3 - r_{23}^3) = 0.$$
(5)

Taking the difference of equations (4) and (5) and factoring the result yields

$$(r_{13}-r_{24})(r_{13}^2r_{14}^2r_{24}^2+r_{13}^2r_{23}^2r_{24}^2+r_{13}r_{14}^3r_{23}^2+r_{13}r_{14}^2r_{23}^3+r_{14}^3r_{23}^2r_{24}+r_{14}^2r_{23}^3r_{24})=0.$$

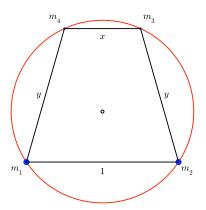


Figure: An example of an isosceles trapezoid central configuration. The center of the circumscribing circle is marked with an O.

We set $m_1 = m_2 = 1$. Then, $m_3 = m_4 = m$ where

$$m = \frac{x^2(1-y^3)}{y^3-x^3}.$$

In order for the isosceles trapezoid to be a c.c., we must have

$$f = (y^2 + x)^{3/2}(2y^3 - x^3 - 1) - y^3 - x^3y^3 + 2x^3 = 0.$$
 (6)

Examining f=0, we see that while the smallest side of the trapezoid (parallel to the base) can range from 0 to 1, the length of the congruent legs is considerably constrained between approximately 0.91 and 1.

Conjecture: y is a differentiable function of x and m is an increasing function of x ranging from 0 to 1.

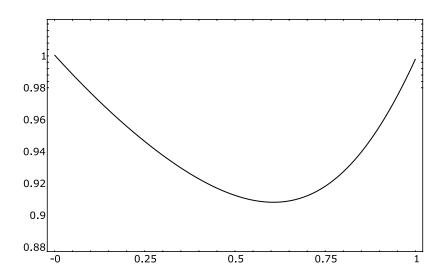


Figure: The relationship between the two distances $x = r_{34}$ and $y = r_{14} = r_{23}$ in the isosceles trapezoid family of cc's.

Theorem

The maximum arclength θ_{12} in the isosceles trapezoid family of cc's is a monotonically decreasing function of the smallest sidelength x. The maximum of θ_{12} is 120° at the equilateral triangle ($m_3=m_4=0$) and $r_{34}=0$) while the minimum is 90° attained at the square configuration.

Note that these bounds are the same as those for the kite family of cyclic central configurations.

Some Geometric Properties of CCC's

- The maximum interior angle of a 4-body cyclic central configuration lies between 90° and 120° while the minimum interior angle lies between 60° and 90°.
- Any cyclic c.c. containing two opposite bodies (1 opposite 3 or 2 opposite 4) on a diameter of the circumscribing circle must be a kite configuration.
- No cyclic c.c. can lie entirely in a semi-circle.
- ullet Three bodies of a c.c.c. cannot lie on the same half of the circle as the longest side r_{12} unless the configuration is a kite.

Some Open Questions/Future Work

For a cyclic quadrilateral ordered sequentially, the lengths of each diagonal r_{13} and r_{24} can be written as

$$r_{13} = \sqrt{\frac{ab}{c}}$$
 and $r_{24} = \sqrt{\frac{ac}{b}}$ (7)

where $a = r_{12}r_{34} + r_{14}r_{23}$, $b = r_{12}r_{14} + r_{23}r_{34}$ and $c = r_{12}r_{23} + r_{14}r_{34}$.

If the two equations in (7) hold, then both P=0 and V=0 follow. Substituting formulae (7) for r_{13} and r_{24} into

$$(r_{13}^3 - r_{12}^3)(r_{23}^3 - r_{34}^3)(r_{24}^3 - r_{14}^3) \ = \ (r_{12}^3 - r_{14}^3)(r_{24}^3 - r_{34}^3)(r_{13}^3 - r_{23}^3)$$

and setting $r_{12} = 1$ yields a complicated equation in three variables,

$$F(r_{14},r_{23},r_{34}) = 0.$$



Let Ω denote the set of mutual distances satisfying

$$r_{13} \geq r_{24} > r_{12} = 1 \geq r_{14} \geq r_{23} \geq r_{34}$$

and $\pi(\Omega)$ its projection onto $r_{14}r_{23}r_{34}$ space.

Let Γ denote the intersection of F=0 and $\pi(\Omega)$. Any point on the surface Γ yields a c.c.c. with positive masses.

Let $G: \Gamma \mapsto \mathbb{R}^3$ be the mass map taking a point in Γ to a 3-tuple of positive mass ratios.

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What do the surfaces F=0 and Γ look like? What does the image of G look like in the mass space? The boundaries are determined by the kite and isosceles trapezoid configurations.